

Exam for “Symmetries and quantum field theory”

Master 2 Concepts fondamentaux de la physique, parcours physique quantique

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Duration : 3h.

Handwritten notes and typed sheets given during the lectures/exercise sessions are the only documents allowed. The two problems are independent. We suggest to spend about 1.5h on each problem.

1 Axion electrodynamics for 3+1 topological insulators

We consider a 3D insulating (i.e. which does not conduct electricity) material featuring a magneto-electric effect. The latter is an unusual effect : for example, an electric field induces a magnetic polarization and a magnetic field produces an electric polarization. The electrodynamics of such an insulator can be shown to be described by the following action¹ :

$$S = S_0 + S_\theta = \int_M d^4x \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{\alpha}{2\pi} \epsilon^{\mu\nu\rho\sigma} \theta (\partial_\mu A_\nu) (\partial_\rho A_\sigma) \right)$$

where M is the space-time manifold, $\alpha = e^2/(4\pi\hbar)$ is the fine structure constant and $\epsilon^{\mu\nu\rho\sigma} \equiv \epsilon_{\mu\nu\rho\sigma}$ is the totally antisymmetric symbol such that $\epsilon^{0123} = +1$ (we remind that it has the same expression in any frame). The first term is the Maxwell action and depends on the gauge field A^μ . The second is the θ -term and depends on a real field $\theta(x)$ called the axion field. This field varies slowly in space-time and plays the role of an order parameter field : it has no dynamics on its own (there is no kinetic energy for this field) and may be thought of as an external field.

We first study the nature of the θ -term **when $\theta = \text{constant}$** (as relevant for a bulk homogeneous crystal).

1. Show that S_0 is a Lorentz scalar. Show that S_θ is a Lorentz pseudo-scalar. The following identity, valid for a 4×4 matrix C with elements C^μ_ν , may be useful : $\epsilon_{\mu\nu\rho\sigma} C^\mu_\alpha C^\nu_\beta C^\rho_\gamma C^\sigma_\delta = \epsilon_{\alpha\beta\gamma\delta} \det(C)$.

2. Write the action S as a function of the electric \vec{E} and magnetic \vec{B} fields. Is S_θ gauge invariant? Write S_θ as a function of the electromagnetic field strength $F^{\mu\nu}$ and its dual $\tilde{F}^{\mu\nu}$.

3. Obtain the Euler-Lagrange equation for the A_μ field. Write the corresponding Maxwell equations for \vec{E} and \vec{B} . What is the effect of the θ term?

4. Starting from the expression of S_θ in terms of the gauge field A_μ show that S_θ can be written as a total derivative. Use this property to express S_θ as an integral over the manifold boundary ∂M . In hindsight, explain the results of 3).

We admit that θ is actually an angular variable and therefore defined modulo 2π . We take $\theta \in [0, 2\pi[$ in the following. The properties obtained in 3) and 4) mean that S_θ is a topological term, i.e. it only depends on the topology of the field A^μ (the way the field wraps around the space-time manifold) and not on its precise configuration.

5. What is the behavior of $\vec{E} \cdot \vec{B}$ under a parity transformation P ? Under a time-reversal transformation T ? Under PT ? From this considerations, deduce the behavior of S_θ and of S_0 under P , T and PT .

6. Assume that the action has time reversal symmetry (TRS). What are the possible values of θ ? A trivial insulator corresponds to $\theta = 0$. Knowing that a topological insulator can be defined as an insulating crystal with TRS and which is distinct from a trivial insulator, what is its θ value?

What are the possible θ values for a 3D parity invariant insulator? And for a 3D insulator without P and T symmetries?

We now take a general situation **in which $\theta \neq \text{constant}$** as relevant for an inhomogeneous crystal, e.g. (remember that θ has no dynamics and can be thought of as a given external field).

7. Obtain the Euler-Lagrange equation for the A_μ field and give the corresponding equations for \vec{E} and \vec{B} when $\theta \neq \text{constant}$. What is the effect of the axion field? Does it affect all four Maxwell equations? These modified Maxwell equations are known as “axion electrodynamics”.

1. Starting from S.I. units, we take $c \equiv 1$ and $\epsilon_0 \equiv 1$ (hence $\mu_0 = 1$) so that the fine structure constant $\alpha \equiv e^2/(4\pi\epsilon_0\hbar c) = e^2/(2\hbar) \approx 1/137$ and the conductance quantum $e^2/h = 2\alpha$.

8. Compute the (topological part of the) electrical current defined by $J^\mu = -\frac{\delta S_\theta}{\delta A_\mu}$ and show that it vanishes if $\theta = \text{constant}$.

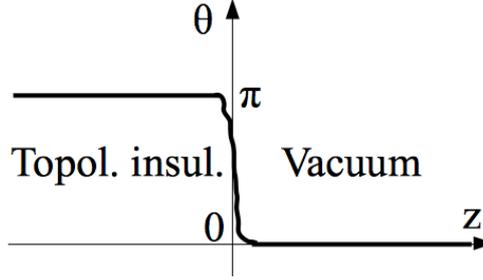


FIGURE 1 – Domain wall in the axion field $\theta(z)$ at the surface of a topological insulator.

9. Consider the surface (at $z = 0$) of a semi-infinite 3D topological insulator ($z < 0$ is the topological insulator and $z > 0$ is the vacuum which is assumed to be a trivial insulator). This is represented as a domain wall in the field $\theta(t, x, y, z) = \theta(z)$, see figure 1. Next apply an electric field $\vec{E} = E_y \vec{e}_y$ (with $E_y(x, y, z, t) = \text{constant}$ and $\vec{B} = 0$) and compute the 2D current density defined by $\vec{j}(x, y) = \int dz \vec{J}(x, y, z)$. Obtain the Hall conductivity σ_{xy} defined by $j_x = \sigma_{xy} E_y$. Show that – in the absence of a magnetic field – there is a Hall effect and give its value in units of the conductance quantum $e^2/h = 2\alpha$.

10. What determines the direction (along \vec{e}_x or $-\vec{e}_x$) of the Hall current? Does this make sense?

References :

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X.L. Qi, T.L. Hugues and S.C. Zhang, Phys. Rev. B **78**, 195424 (2008).

A.M. Essin, J.E. Moore and D. Vanderbilt, Phys. Rev. Lett. **102**, 146805 (2009).

2 Boron nitride and the 2+1 Dirac equation

We consider the low energy properties of conduction electrons in boron nitride (BN). This 3D material is made of stacked 2D planes. Each plane is a honeycomb lattice (see figure 2(a)), with nearest neighbor distance $a \sim 1\text{\AA}$. In the following, we study a single layer of BN and neglect the spin of the electrons. In a tight-binding model, each atom (either B or N) may accommodate a single conduction electron. The latter can hop from an atom to a nearest neighbor atom with a probability amplitude $t \sim 1$ eV. The on-site energy ϵ_α for an electron depends on the atom type ($\alpha = N$ or B). We choose the zero of energy such that $\epsilon_N = -\epsilon_B > 0$. We call $c_{i\alpha}^\dagger$ the creation operator of an electron on the atom (i, α) , where i indicates a triangular Bravais lattice site and α which of the two atoms in the unit cell (see figure 2(a)). The Hamiltonian is given by :

$$H = \sum_{i,\alpha} \epsilon_\alpha c_{i\alpha}^\dagger c_{i\alpha} - t \sum_{\langle i\alpha, j\alpha' \rangle} \left(c_{i\alpha}^\dagger c_{j\alpha'} + \text{H.c.} \right) \quad (1)$$

where the notation $\langle i\alpha, j\alpha' \rangle$ indicates that the sum is over every pair of nearest neighbor atoms in the lattice.

1. Performing a Fourier transform on the Bravais lattice index, show that the Hamiltonian may be written as :

$$H = \sum_{\vec{k} \in \text{BZ}} \begin{pmatrix} c_{\vec{k}N}^\dagger & c_{\vec{k}B}^\dagger \end{pmatrix} \cdot \begin{pmatrix} \Delta & f(\vec{k})^* \\ f(\vec{k}) & -\Delta \end{pmatrix} \cdot \begin{pmatrix} c_{\vec{k}N} \\ c_{\vec{k}B} \end{pmatrix} = \sum_{\vec{k} \in \text{BZ}} \Psi_{\vec{k}}^\dagger \cdot H(\vec{k}) \cdot \Psi_{\vec{k}} \quad (2)$$

where BZ indicates the Brillouin zone (see figure 2(b)) and $\Psi_{\vec{k}}$ is a two-component object for each \vec{k} . Give the expression of the function $f(\vec{k})$ and of Δ .

2. What are the eigen-energies $\epsilon_{\vec{k}}^\pm = \pm \epsilon_{\vec{k}}$ (with $\epsilon_{\vec{k}} > 0$) of the single electron states of (2) ?

3. Show that $\epsilon_{\vec{k}}$ has a minimum at two inequivalent values of \vec{k} on the edges of the BZ : $\vec{k} = \pm \vec{K} = \pm(\frac{4\pi}{3\sqrt{3}a}, 0)$. These two points in the BZ are known as two valleys (K and K').

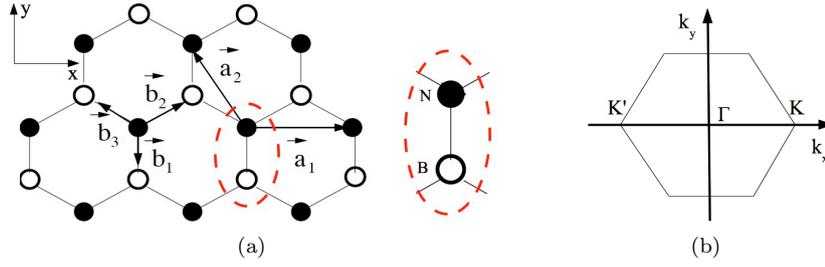


FIGURE 2 – (a) The honeycomb lattice is a triangular Bravais lattice (with lattice vectors \vec{a}_1, \vec{a}_2) with two atoms (N and B) in the unit cell (the two-atom basis is indicated by a dashed ellipse). Nearest neighbor vectors \vec{b}_1, \vec{b}_2 and \vec{b}_3 are also defined. (b) The corresponding hexagonal Brillouin zone (remarkable points Γ, K and K' are indicated).

4. By expanding the single electron Hamiltonian $H(\vec{k})$ at low energy in the vicinity of $+\vec{K}$ and $-\vec{K}$ (with the notation $\vec{k} = \vec{p} + \xi\vec{K}$, $\xi = \pm$ being the valley index), show that the effective theory becomes (with $\hbar = 1$)

$$H(\vec{k}) \approx \sum_{\xi=\pm} H(\vec{k} = \vec{p} + \xi\vec{K}) \approx \sum_{\xi=\pm} H_{\xi}(\vec{p})$$

$$\text{with } H_{\xi} = v_F \vec{\alpha}_{\xi} \cdot \vec{p} + mv_F^2 \beta \quad (3)$$

where $\vec{\alpha}_{\xi} = (\xi\sigma_x, -\sigma_y)$ and $\beta = \sigma_z$ in terms of the Pauli matrices σ_x, σ_y and σ_z . Give the expression of v_F and mv_F^2 as a function of the tight-binding model parameters. Estimate v_F and give an order of magnitude for m both in S.I. units.

In the following, we use units such that $v_F = 1$ in addition to $\hbar = 1$.

5. The single electron Hamiltonian (3) may be rewritten in position representation by identifying \vec{p} with $-i\vec{\nabla}$. Show that this gives a Dirac Hamiltonian for each valley $\xi = \pm$.

6. In order to write the Dirac equation following from (3), introduce the corresponding 2+1 dimensional γ^{μ} matrices (How many such matrices? What is their algebra? Their minimal size? In the representation of the above questions, give their expression in terms of Pauli matrices). What is the physical meaning of the components of the Dirac spinor $\Psi(x)$ here?

7. By analogy with the situation in 3+1 dimension, we define γ^5 by $\gamma^5 = \frac{i}{3!}\epsilon_{\mu\nu\rho}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}$. Compute this matrix and its anticommutator with γ^{μ} . Does it define a chirality transformation which differs from a $U(1)$ phase transformation? Conclude on the existence of chirality in 2+1 dimension.

8. The system is now placed in a constant and uniform magnetic field \vec{B} perpendicular to the BN layer. How is the Dirac Hamiltonian (3) affected by the magnetic field? Write also the corresponding Dirac equation.

9. We choose a gauge such that $A^{\mu} = (A^0, A^1, A^2) = (0, 0, Bx)$ with $B > 0$. Using the fact that the Hamiltonian commutes with p_y , obtain its spectrum and show that the energy of the n^{th} Landau level is given by :

$$\epsilon_n = \pm\sqrt{m^2 + 2neB} \text{ if } n = 1, 2, 3, \dots \text{ and } \epsilon_0 = -m\xi \text{ if } n = 0 \quad (4)$$

References :

- G.W. Semenoff, Phys. Rev. Lett. **53**, 2449 (1984).
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